

On random permutations with nonuniform distributions

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Program

- Quantum Bose gas

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- Model of spatial permutations (“general”)

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- Illustration: lattice permutations
- Results for spatial permutations with small cycle weights

From the Bose gas to spatial permutations

n bosons in $\Lambda \subset \mathbb{R}^d$

State space $L^2_{\text{sym}}(\Lambda^n)$

Hamiltonian

$$H = -\sum_{i=1}^n \Delta_i + \sum_{i < j} U(x_i - x_j)$$

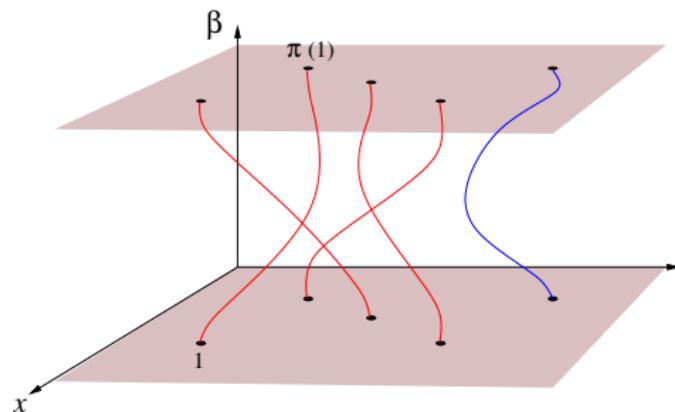
Feynman-Kac formula for

$$Z = \text{Tr} e^{-\beta H} :$$

with $\mathbf{x} = (x_1, \dots, x_n)$,

$$Z = \int_{\Lambda^n} d\mathbf{x} \sum_{\pi \in \mathcal{S}_n} e^{-H(\mathbf{x}, \pi)} \quad \text{with "Hamiltonian"} \quad H : \Lambda^n \times \mathcal{S}_n \rightarrow \mathbb{R}$$

(Explicit expression for H in terms of Wiener trajectories)



Model of spatial permutations (general)

State space: $\Omega_{\Lambda,n} = \Lambda^n \times \mathcal{S}_n$

Gibbs weight $e^{-H(\mathbf{x},\pi)}$

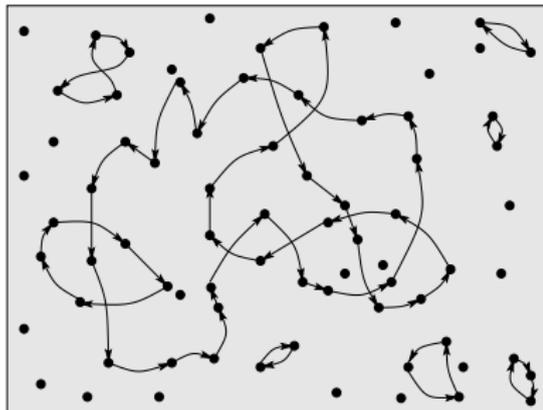
Expectation of random variable Θ :

$$E_{\Lambda,n}(\Theta) = \frac{1}{Z} \int_{\Lambda^n} d\mathbf{x} \sum_{\pi} \Theta(\mathbf{x}, \pi) e^{-H(\mathbf{x},\pi)}$$

General Hamiltonian:

$$H(\mathbf{x}, \pi) = \sum_{i=1}^n \xi(x_i - x_{\pi(i)}) + \sum_{i < j} V(x_i, x_{\pi(i)}; x_j, x_{\pi(j)})$$

$V(x_i, x_{\pi(i)}; x_j, x_{\pi(j)})$ represents the pair interaction between the jumps $x_i \mapsto x_{\pi(i)}$ and $x_j \mapsto x_{\pi(j)}$



Model of spatial permutations (with cycle weights)

Bose gas: $\xi(x) = \frac{1}{4\beta}|x|^2$; interactions between $x \mapsto y$ and $x' \mapsto y'$ given by

$$\begin{aligned} V(x, y; x', y') &= \int [1 - e^{-\frac{1}{4} \int_0^\beta U(\omega(s)) ds}] dW_{x-x', y-y'}(\omega) \\ &= K(x - x', y - y') \end{aligned}$$

with K the integral kernel of $e^{2\beta\Delta} - e^{2\beta(\Delta-U)}$

After additional calculations and some (exact?) approximations:
Model with cycle weights

$$H(\mathbf{x}, \pi) = \sum_{i=1}^n \xi(x_i - x_{\pi(i)}) + \sum_{j \geq 1} \alpha_j r_j(\pi)$$

with α_j : parameters, and $r_j(\pi)$: number of cycles of length j

Preliminary study: *random permutations with cycle weights*

Random permutations with cycle weights

Probability of permutation $\pi \in \mathcal{S}_n$:

$$P(\pi) = \frac{1}{n!h_n} \exp\left\{-\sum_{j \geq 1} \alpha_j r_j(\pi)\right\}$$

with α_j : parameters, and $r_j(\pi)$: number of cycles of length j

Uniform permutations: $\alpha_j \equiv 0$

Further motivation:

- *Ewens distribution* in mathematical biology: $e^{-\alpha_j} \equiv \theta$. Model for genetic mutations in haploid populations; $\sum r_j$ describes the number of alleles; the Ewens distribution characterizes the stationary state
- This distribution also appears in Bayesian statistics, and in the Pólya urn model
- Links with random partitions

Random partitions

A partition λ of n : $\lambda = (\lambda_1, \lambda_2, \dots, \lambda_k)$ with

$$\lambda_1 \geq \lambda_2 \geq \dots \geq \lambda_k > 0 \quad \text{and} \quad \sum_{j=1}^k \lambda_j = n$$

The length of the partition, k , is not fixed

Occupation numbers: $r_j(\lambda) =$ number of indices i such that $\lambda_i = j$.

Occupation numbers determine the partition, and conversely. They satisfy

$$\sum_{j \geq 1} j r_j = n$$

We consider only random variables that depend on the cycle lengths (r_j) of permutations. Thus the model is really a model of random partitions. The Ewens distribution is often written

$$p(\lambda) = \frac{n!}{(\theta)_n} \prod_{j \geq 1} \frac{\theta^{r_j}}{j^{r_j} r_j!}$$

Literature

Several studies for the number of cycles $K(\pi)$

- **Gontcharov** (1944): central limit theorem for K , for uniform random permutations:

$$\frac{K - \log n}{\sqrt{\log n}} \rightarrow N(0, 1) \quad \text{in distribution}$$

- An amazing generalization of the central limit theorem above: Let $K(t, \pi) = \sum_{j=1}^{n^t} r_j$, and

$$Y(t, \pi) = \frac{K(t, \pi) - \theta t \log n}{\sqrt{\theta \log n}}$$

Then $Y(t, \pi)$ converges weakly to the Wiener measure on the space of càdlàg functions on $[0,1]$. **DeLaurentis & Bittel** (1985) for uniform distributions, **Hansen** (1990) for Ewens

- **Feng & Hoppe** (1998): Large deviations for $\frac{K(t, \pi)}{\log n}$

Results for permutations with cycle weights

We are interested in the distribution of cycle lengths, rather than in the number of cycles. $\ell_1(\pi)$: length of cycle that contains 1 .

Weights $\theta_j = e^{-\alpha_j}$ with $\alpha_j \asymp j^\gamma$

Theorem (with Volker Betz & Yvan Velenik, 2009)

- The case $\gamma < 0$ is close to uniform distribution, $P(\ell < sn) \rightarrow s$
- In the case $\gamma = 0$, i.e. $e^{-\alpha_j} \approx \theta$ (Ewens), $P(\ell_1 > sn) \rightarrow (1 - s)^\theta$ for any $0 \leq s \leq 1$ (cycles are macroscopic)
- The case $0 < \gamma < 1$ is surprising, with $P(\ell_1 > n - k) \geq 1 - o(k)$ as $n \rightarrow \infty$. Almost all indices belong to a single giant cycle!
In addition: (i) $P(\ell_1 = n) > 0$; (ii) in some cases, $P(\ell_1 = 1) > 0$
- The case $\gamma = 1$ corresponds to uniform permutations
- When $\gamma > 1$ cycles become shorter, and $\ell_1 \approx (\frac{1}{\gamma-1} \log n)^{1/\gamma}$ with probability 1

Sketch of proof for small weights

The case $\alpha_j \equiv 0$ is easy:

$$\begin{aligned} P(\ell_1(\pi) \in \{1, \dots, sn\}) &= \sum_{j=1}^{sn} P(\ell_1(\pi) = j) \\ &= \sum_{j=1}^{sn} \frac{1}{n!} (n-1) \dots (n-j+1) \cdot (n-j)! = \frac{sn}{n} = s \end{aligned}$$

Main ideas for the case $\alpha_j \neq 0$, small: ($P(\pi) = \frac{1}{n! h_n} \prod_j e^{-\alpha_j r_j(\pi)}$)

Observe that

$$P(\ell_1 \in [a, b]) = \frac{1}{n} \sum_{j=a}^b \underbrace{e^{-\alpha_j}}_{\approx 1} \underbrace{\frac{h_{n-j}}{h_n}}_{\text{show that } \approx 1}$$

Expression for h_n using generating functions

$$h_n = \frac{1}{n} \sum_{j=1}^n e^{-\alpha_j} h_{n-j} \quad (\text{because } P(\ell_1 \in [1, n]) = 1)$$

Then

$$(n+1)h_{n+1} = \sum_{j=0}^n b_j h_{n-j} \quad (b_j = e^{-\alpha_{j+1}})$$

$$((n+1)h_{n+1}) = b * h$$

Introduce generating functions:

$$G_h(s) = \sum_{n \geq 0} h_n s^n, \quad G_b(s) = \sum_{n \geq 0} b_n s^n$$

Then

$$\begin{aligned} G'_h(s) &= G_b(s)G_h(s), & G_h(0) &= 1 \\ \rightarrow G_h(s) &= \exp \int_0^s G_b(t) dt = \exp \sum_{n \geq 1} \frac{b_{n-1}}{n} s^n \end{aligned}$$

Expanding the exponential, and matching the terms with the generating function, we get

$$h_n = \sum_{k=1}^n \frac{1}{k!} \sum_{\substack{j_1, \dots, j_k \\ j_1 + \dots + j_k = n}} \prod_{i=1}^k \frac{e^{-\alpha_{j_i}}}{j_i}$$

Remark: in the case $\alpha_j \equiv 0$, we get a puzzling identity:

$$\sum_{k=1}^n \frac{1}{k!} \sum_{\substack{j_1, \dots, j_k \\ j_1 + \dots + j_k = n}} \frac{1}{j_1 \dots j_k} = 1$$

for all n !

One can also obtain

$$h_n = \sum_{k=0}^n \frac{1}{k!} \sum_{\substack{j_1, \dots, j_k \geq 1 \\ j_1 + \dots + j_k \leq n}} \prod_{i=1}^k \frac{e^{-\alpha_{j_i}} - 1}{j_i}$$

Then if $\sum \frac{1}{j} |1 - e^{-\alpha_j}| < \infty$, we have

$$h_\infty = \lim_{n \rightarrow \infty} h_n = \exp \sum_{j \geq 1} \frac{e^{-\alpha_j} - 1}{j}$$

and $\frac{h_{n-j}}{h_n} \approx 1$ when $n - j \gg 1$ and $n \gg 1$. The claim of the theorem easily follows.

Lattice permutations

State space: set of permutations of \mathbb{Z}^d (i.e. bijections $\mathbb{Z}^d \rightarrow \mathbb{Z}^d$)
 σ -algebra generated by “cylinder sets” $B_{x,y} = \{\pi : \pi(x) = y\}$

$$P(B_{x,y}) = \lim_{\Lambda \nearrow \mathbb{Z}^d} \frac{1}{Z(\Lambda)} \sum_{\substack{\pi: \pi(x)=y \\ \pi(z)=z \forall z \notin \Lambda}} e^{-H(\pi)}$$

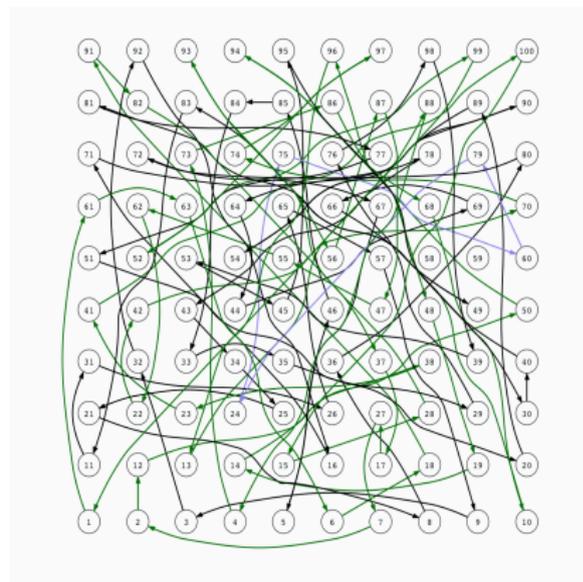
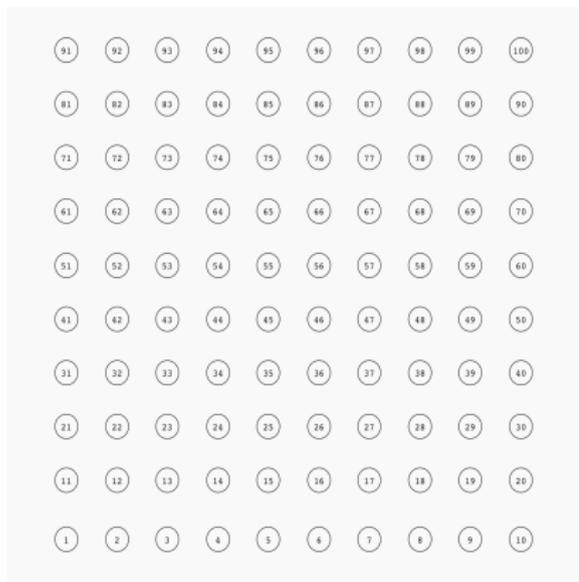
with “Hamiltonian”

$$H(\pi) = \frac{1}{4\beta} \sum_{x \in \Lambda} |x - \pi(x)|^2$$

Existence of limit $\Lambda \nearrow \mathbb{Z}^d$: not clear, requires compactness argument.
Extension to a measure? Yes, provided $\sum_y P(B_{x,y}) = 1$ for all x , but this property has not been proved yet.

Monte Carlo simulations (dimension $d = 2$)

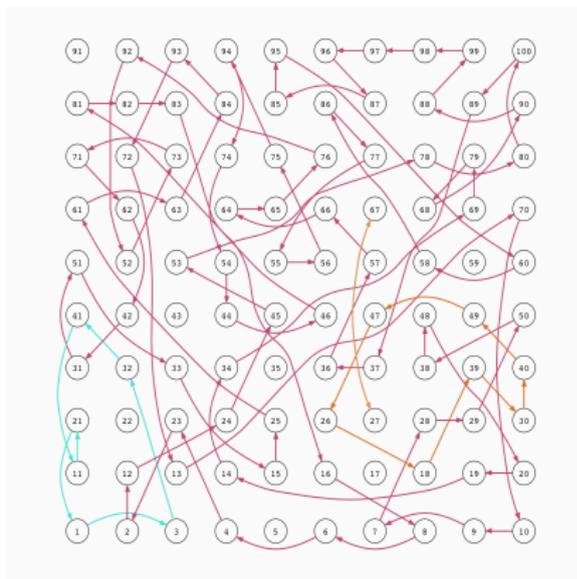
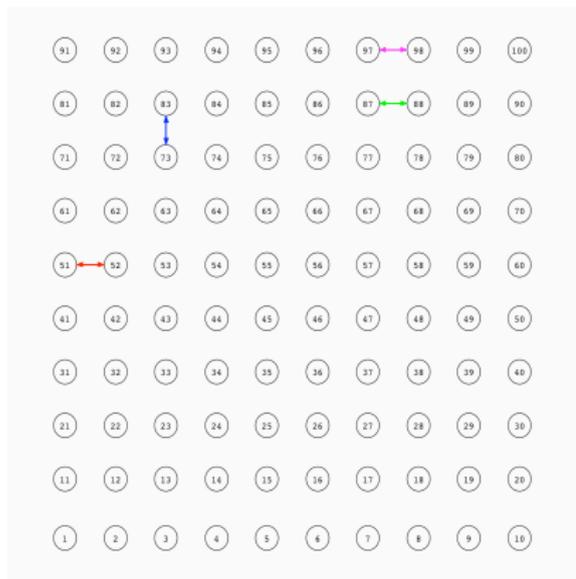
Step 0



(Courtesy of Daniel Gandolfo and Jean Ruiz, CNRS, Marseille)

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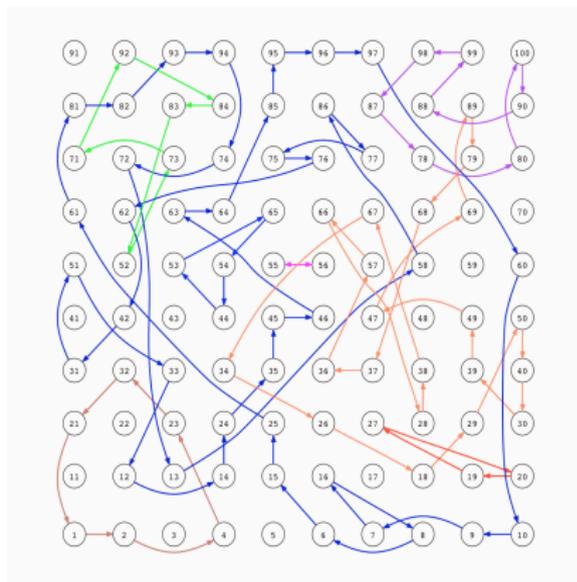
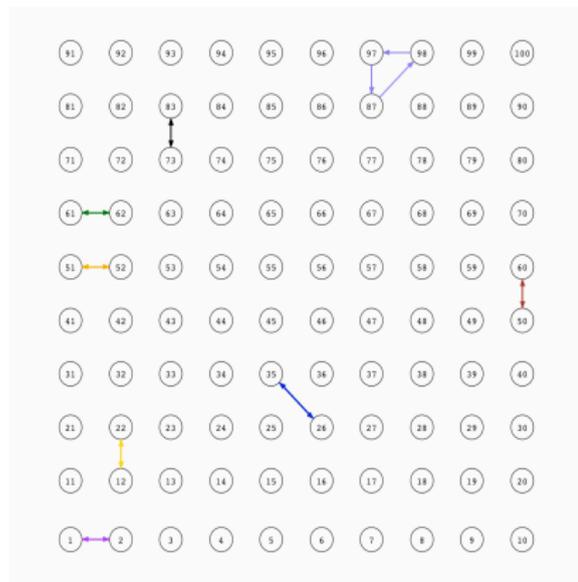
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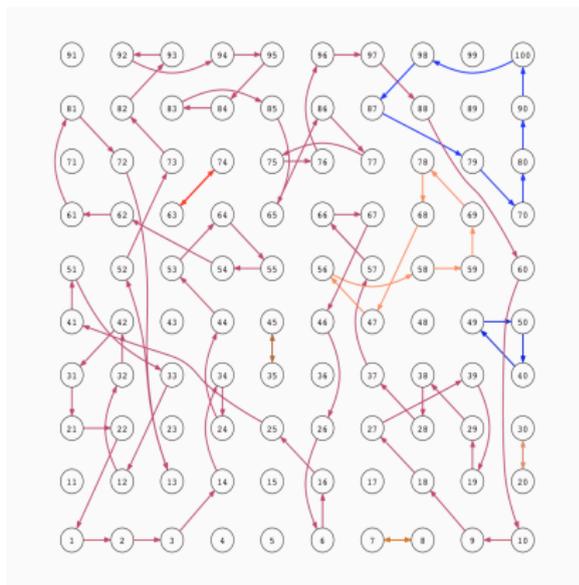
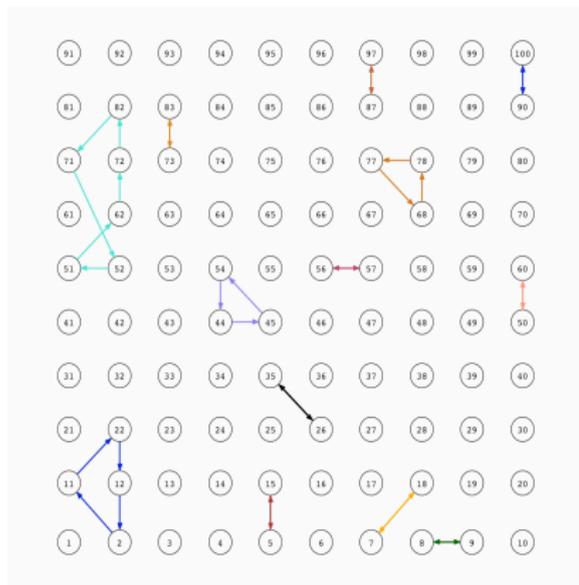
Step 2



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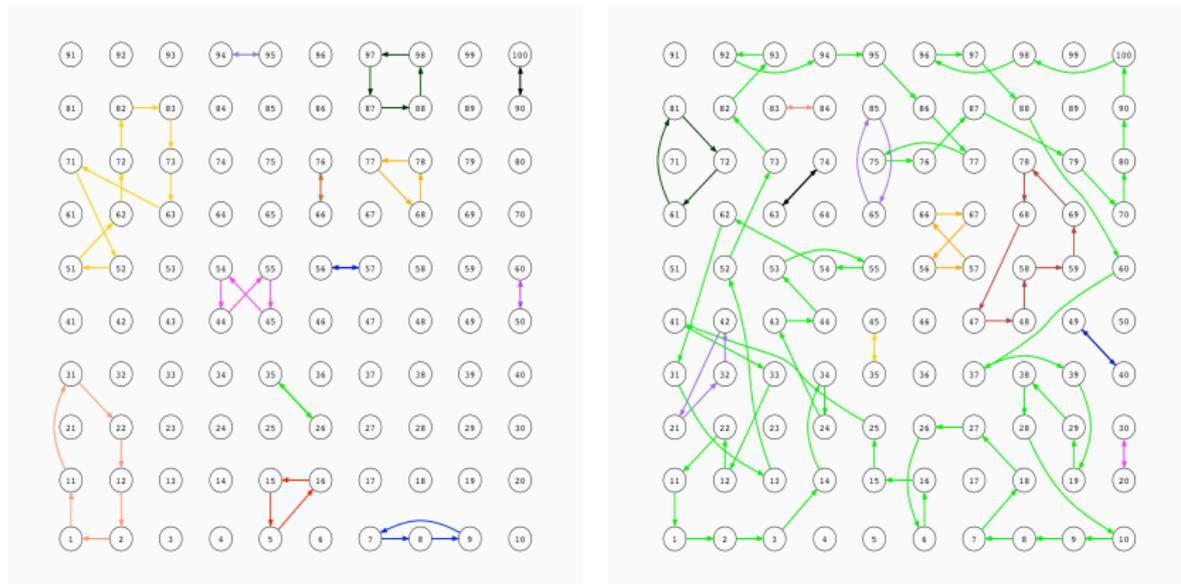
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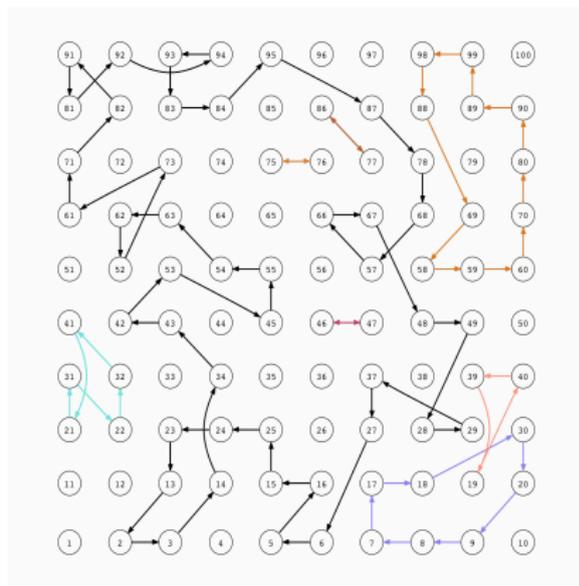
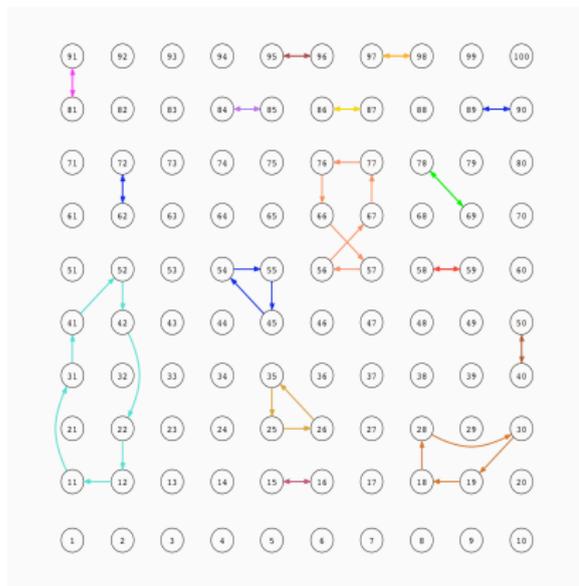
Step 5



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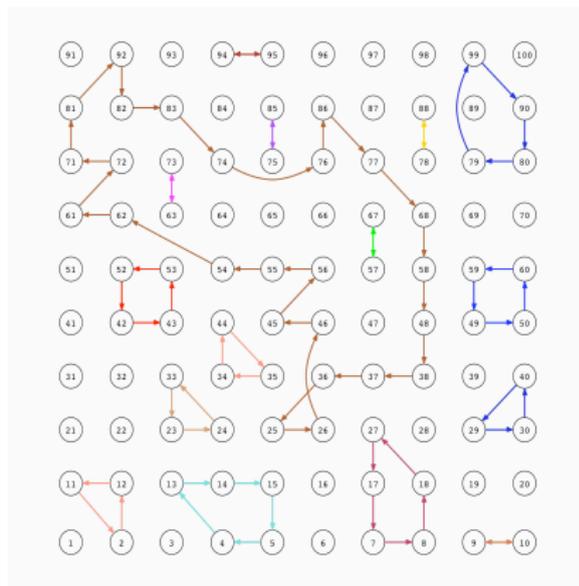
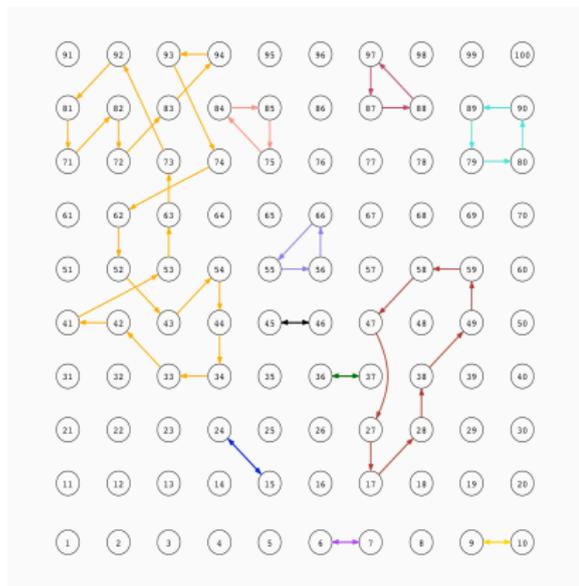
Step 10



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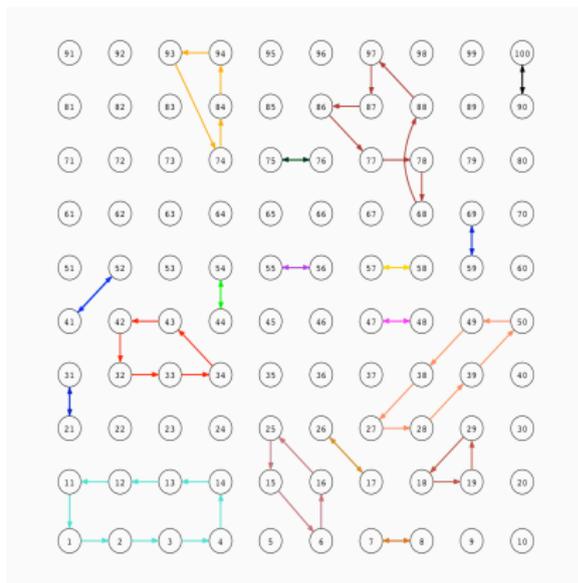
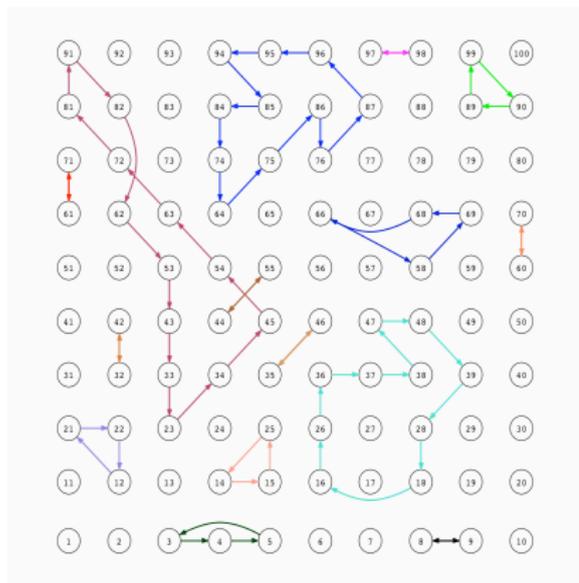
Step 30



(Courtesy of Daniel Gandolfo and Jean Ruiz, CNRS, Marseille)

Monte Carlo simulations (dimension $d = 2$)

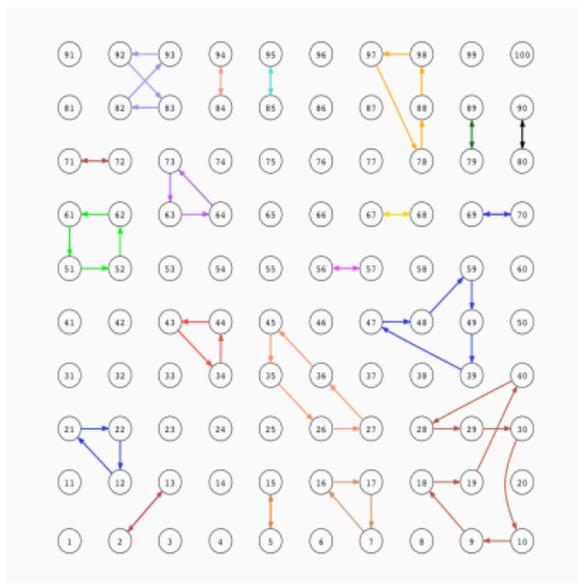
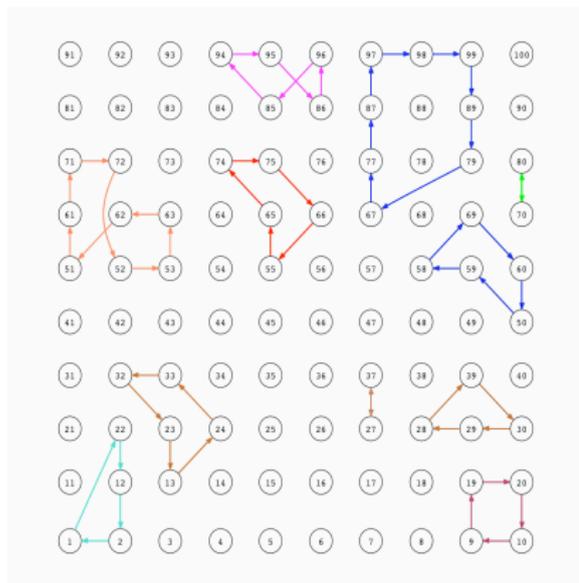
Step 100



(Courtesy of Daniel Gandolfo and Jean Ruiz, CNRS, Marseille)

Monte Carlo simulations (dimension $d = 2$)

Step ∞



(Courtesy of Daniel Gandolfo and Jean Ruiz, CNRS, Marseille)

Monte Carlo simulations: cycle distribution

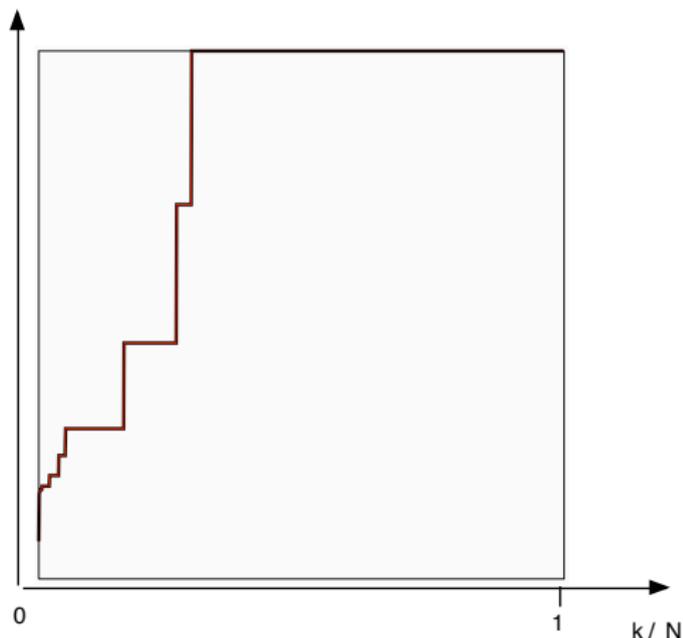
Simulation of a large $3d$ system in equilibrium.

Cubic lattice with $L = 30$, $\beta = 0.5$.

The graph shows the number of sites in cycles of length $\leq k$, as a function of k/N .

red = current

black = average



Gandolfo, Ruiz, U (JSP '07)

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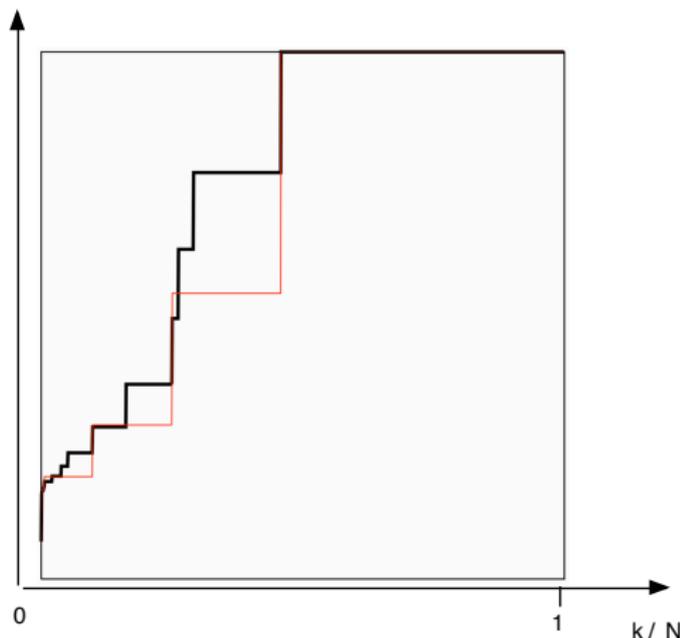
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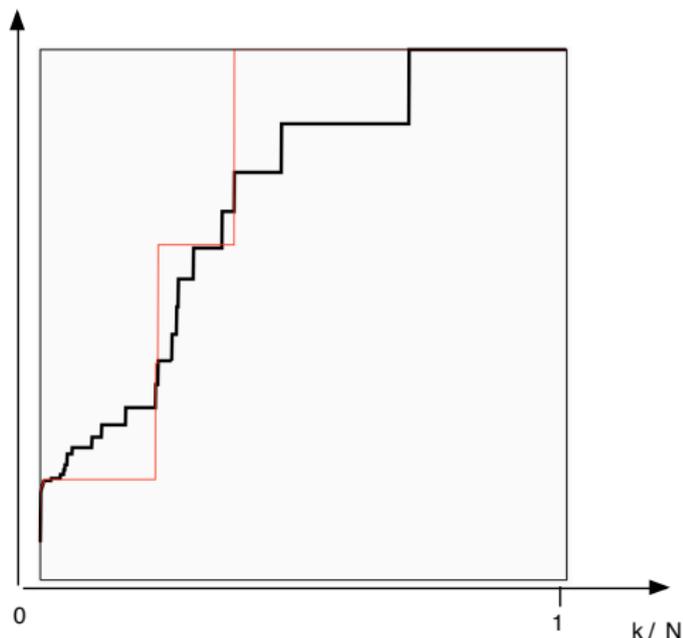
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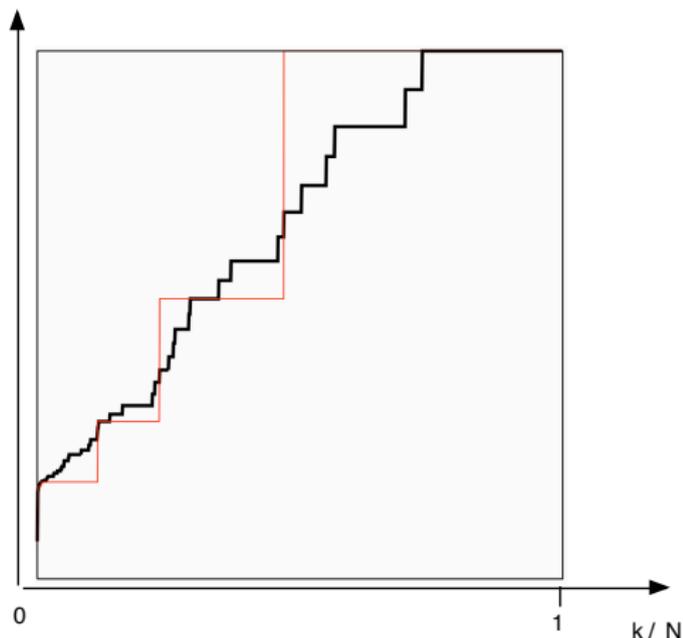
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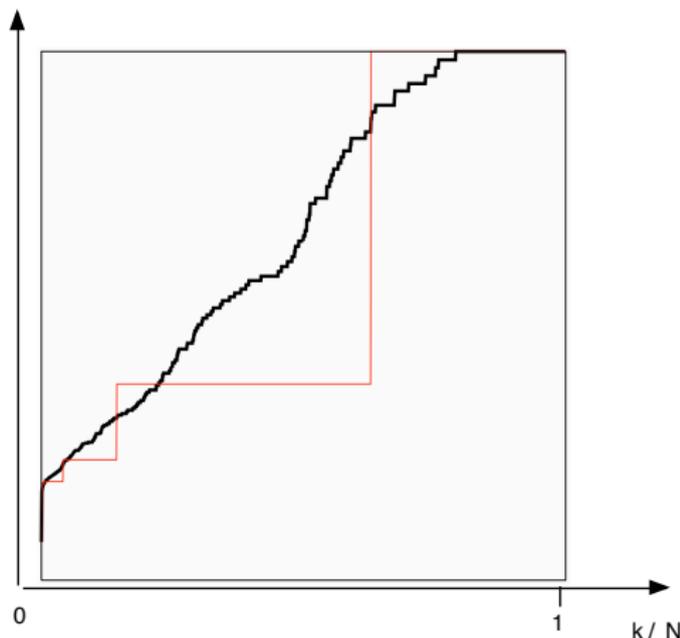
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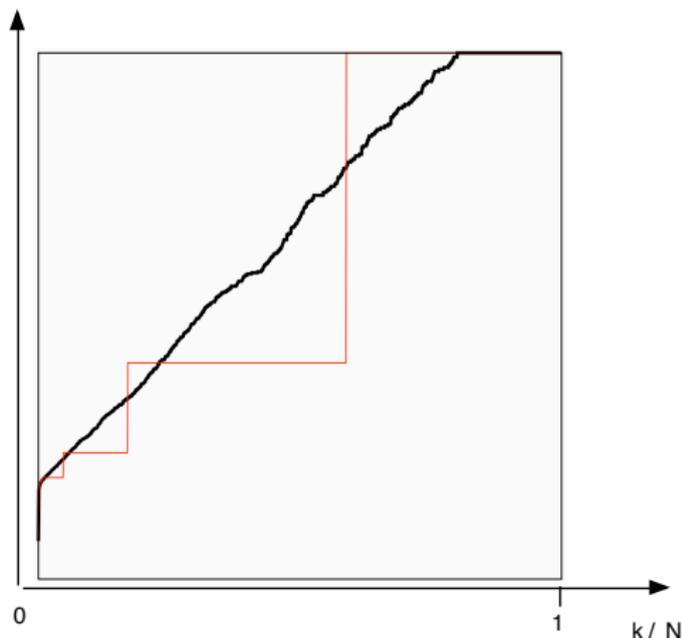
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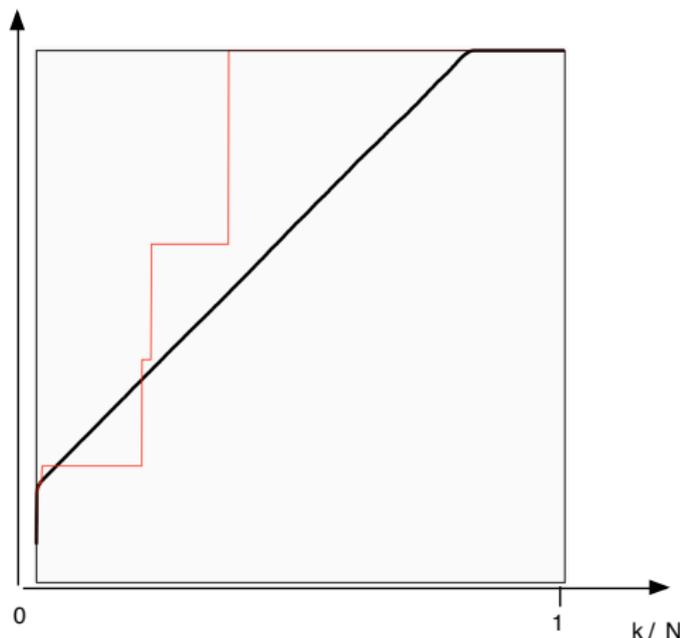
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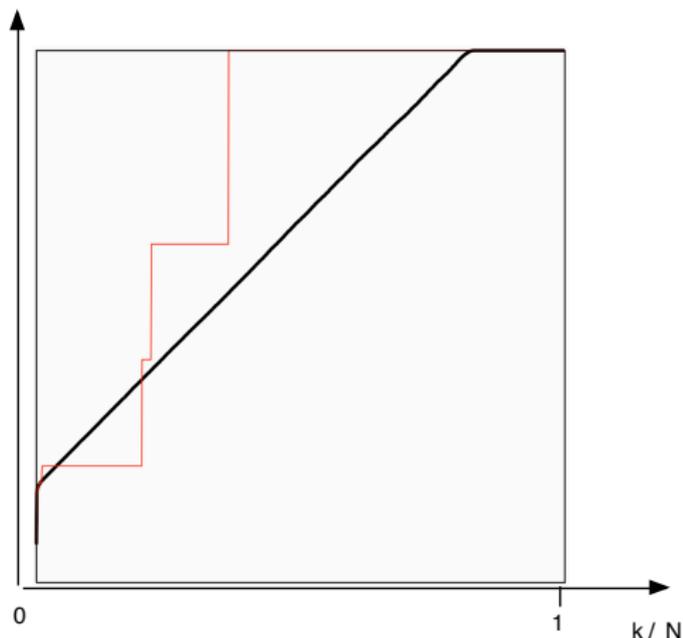
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Conclusion: “infinite” cycles appear!



Gandolfo, Ruiz, U (JSP '07)

Results for spatial permutations with cycle weights

“Annealed” model, with $E_{\Lambda,n}(\Theta) = \frac{1}{Z} \int_{\Lambda^n} d\mathbf{x} \sum_{\pi} \Theta(\mathbf{x}, \pi) e^{-H(\mathbf{x}, \pi)}$ and $H(\mathbf{x}, \pi) = \sum_{i=1}^n \xi(x_i - x_{\pi(i)}) + \sum_{j \geq 1} \alpha_j r_j(\pi)$

We suppose that $\xi(x) = \xi(|x|)$ is increasing, $\int e^{-\xi} = 1$, $\widehat{e^{-\xi}} \geq 0$, and that $\alpha_j \rightarrow 0$ faster than $1/\log j$

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Theorem (with Volker Betz, 2008)

Assume that $\eta(V) \rightarrow \infty$ with $\eta(V)/V \rightarrow 0$, and that $s \geq 0$. Then

- (a) $\lim_{V \rightarrow \infty} P_{\Lambda, \rho V}(\ell_1 \in [1, \eta(V)]) = \begin{cases} 1 & \text{if } \rho \leq \rho_c \\ \rho_c / \rho & \text{if } \rho \geq \rho_c \end{cases}$
- (b) $\lim_{V \rightarrow \infty} P_{\Lambda, \rho V}(\ell_1 \in [\eta(V), \frac{V}{\eta(V)}]) = 0$
- (c) $\lim_{V \rightarrow \infty} P_{\Lambda, \rho V}(\ell_1 \in [\frac{V}{\eta(V)}, sV]) = \begin{cases} 0 & \text{if } \rho \leq \rho_c \\ s/\rho & \text{if } 0 \leq s \leq \rho - \rho_c \\ 1 - \rho_c/\rho & \text{if } 0 \leq \rho - \rho_c \leq s \end{cases}$

Critical density

The critical density is exactly known!

Define $\varepsilon(k)$ by

$$e^{-\varepsilon(k)} = \int_{\mathbb{R}^d} e^{-2\pi i k x} e^{-\xi(x)} dx$$

Notice that $\varepsilon(0) = 0$ and $\varepsilon(k) > a|k|^2$ near $k = 0$. Then

$$\rho_c = \sum_{j \geq 1} e^{-\alpha_j} \int_{\mathbb{R}^d} e^{-j\varepsilon(k)} dk$$

Notice that infinite cycles also occur in $d = 1$ with an integrable jump function, e.g.

$$e^{-\xi(x)} = (|x| + 1)^{-\gamma}$$

with $1 < \gamma < 2$. Then $\varepsilon(k) \sim |k|^{\gamma-1}$ near 0, and

$$\rho_c = \int_{\mathbb{R}^d} \frac{dk}{e^{\varepsilon(k)} - 1} < \infty$$

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- **Many open questions!** Suitable for analytical or numerical studies

THANK YOU!