

Nonequilibrium variational principles from dynamical fluctuations

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To be discussed

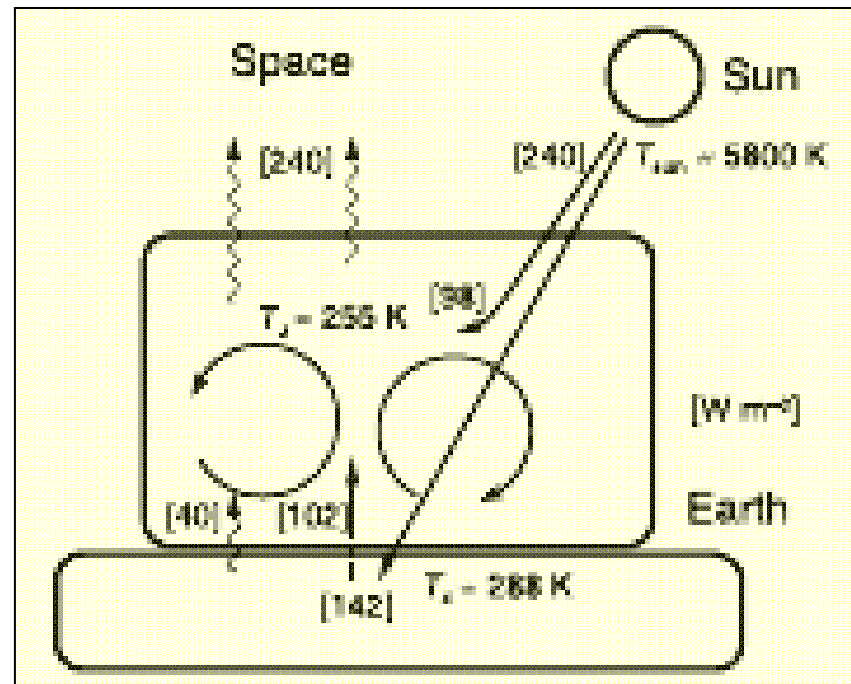
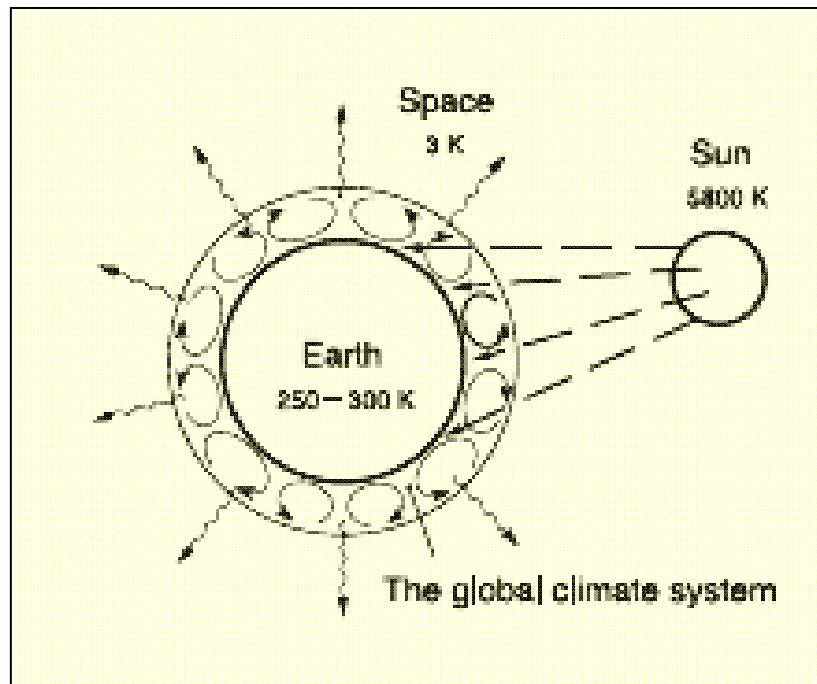
- **Min- and Max-entropy production** principles: various examples
- From variational principles to **fluctuation laws**: equilibrium case
- **Static** versus **dynamical** fluctuations
- **Onsager-Machlup** equilibrium dynamical fluctuation theory
- **Stochastic** models of nonequilibrium
- Conclusions, open problems, outlook,...

❖ *In collaboration with **C. Maes**,
B. Wynants, and **S. Bruers**
(K.U.Leuven, Belgium)*



Motivation: Modeling Earth climate

[Ozawa et al, *Rev. Geoph.* **41** (2003) 1018]



$$\dot{S}_{\text{whole (univ)}} \approx \dot{S}_{\text{surr}} = \left(\frac{1}{T_a} - \frac{1}{T_{\text{sun}}} \right) 240 \approx 0.90 \text{ (W K}^{-1} \text{ m}^{-2}) = \dot{S}_{\text{turb}} + \dot{S}_{\text{abs (short,s)}} + \dot{S}_{\text{abs (short,a)}} + \dot{S}_{\text{abs (long,a)}}$$

$$= \left(\frac{1}{T_a} - \frac{1}{T_s} \right) 102 + \left(\frac{1}{T_s} - \frac{1}{T_{\text{sun}}} \right) 142 + \left(\frac{1}{T_a} - \frac{1}{T_{\text{sun}}} \right) 98 + \left(\frac{1}{T_a} - \frac{1}{T_s} \right) 40$$

Linear electrical networks

explaining MinEP/MaxEP principles

- Kirchhoff's loop law:

$$\sum_k U_{jk} = \sum_k E_{jk}$$

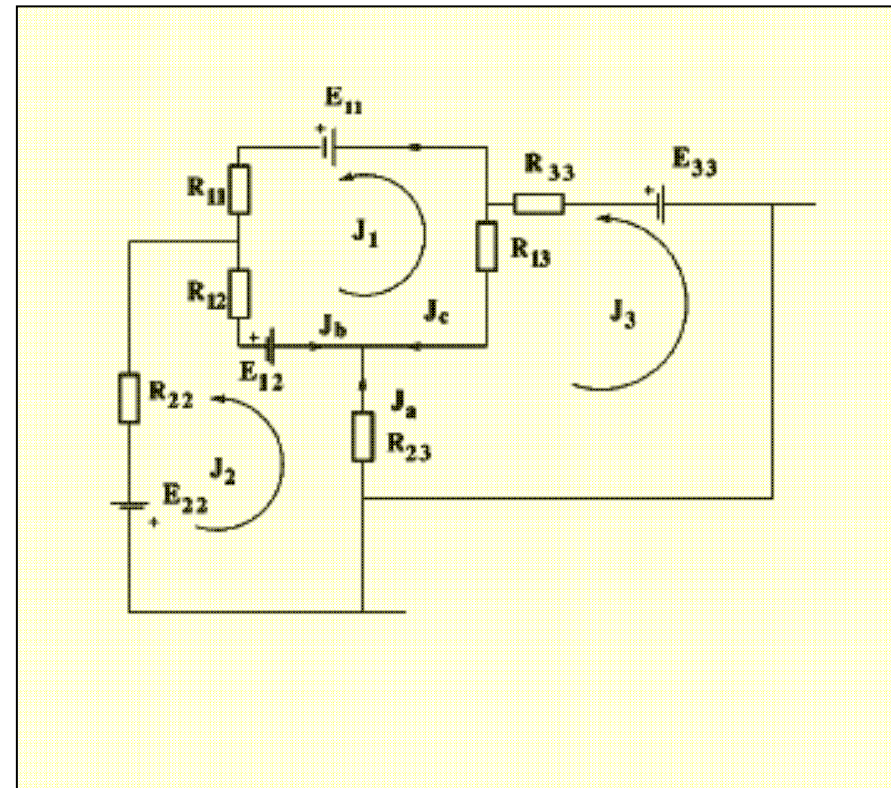
- Entropy production rate:

$$\sigma(U) = \beta Q(U) = \beta \sum_{j,k} \frac{U_{jk}^2}{R_{jk}}$$

- **MinEP** principle:

Stationary values of voltages minimize the entropy production rate

- **Not valid under inhomogeneous temperature!**



Linear electrical networks

explaining MinEP/MaxEP principles

- Kirchhoff's current law:

$$\sum_j J_{jk} = 0$$

- Entropy production rate:

$$\sigma(J) = \beta Q(J) = \beta \sum_{j,k} R_{jk} J_{jk}^2$$

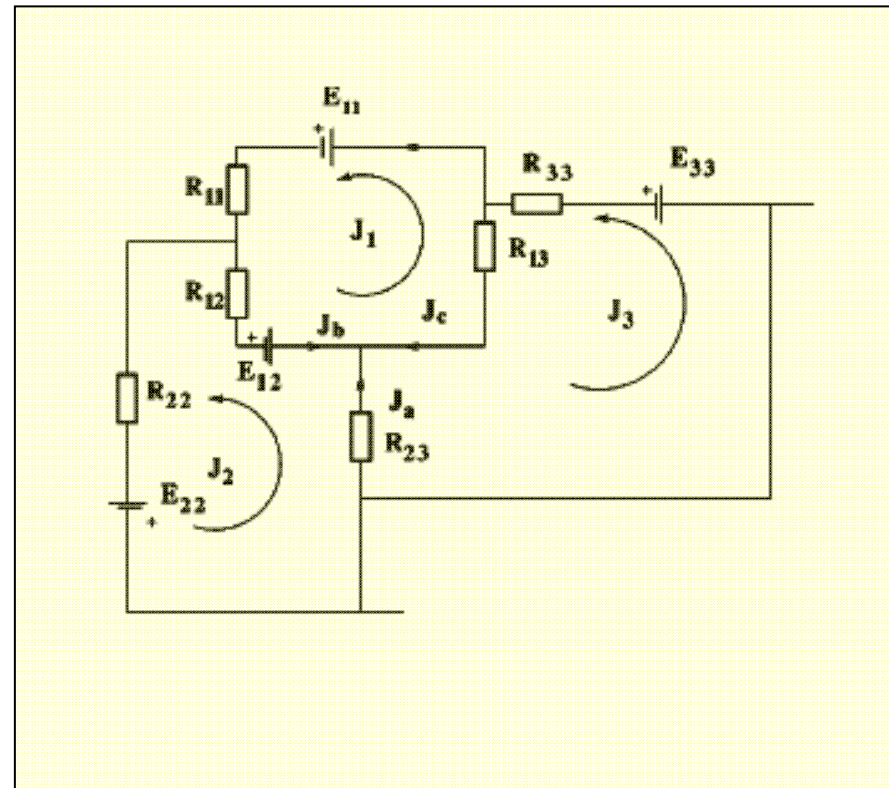
- Work done by sources:

$$W(J) = \sum_{jk} E_{jk} J_{jk}$$

- (Constrained) **MaxEP** principle:

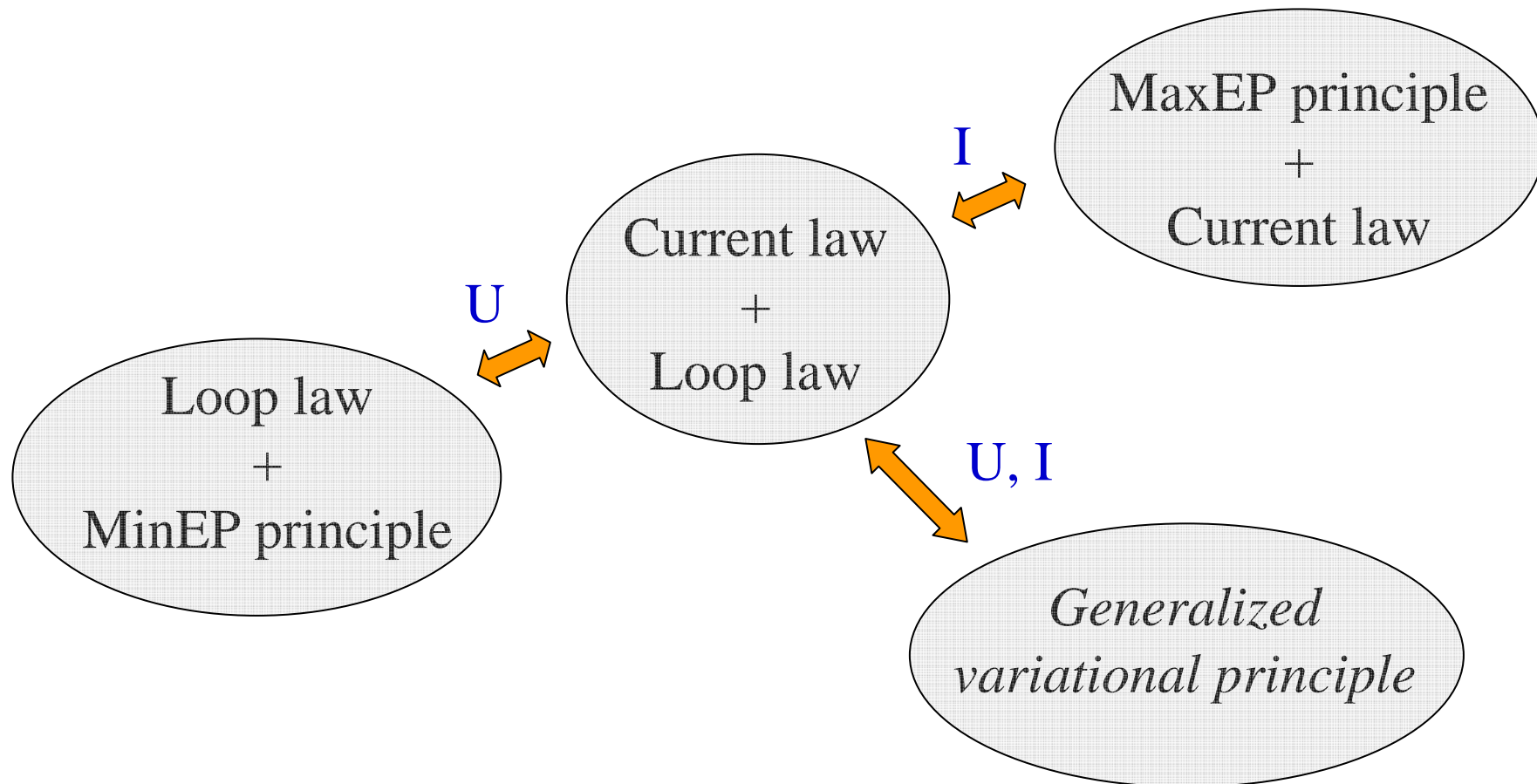
*Stationary values of currents
maximize the entropy
production under constraint*

$$Q(J) = W(J)$$



Linear electrical networks

summary of MinEP/MaxEP principles



From principles to fluctuation laws

Questions and ideas

- How to go beyond **approximate** and *ad hoc* thermodynamic principles?

- Inspiration from thermostatics:

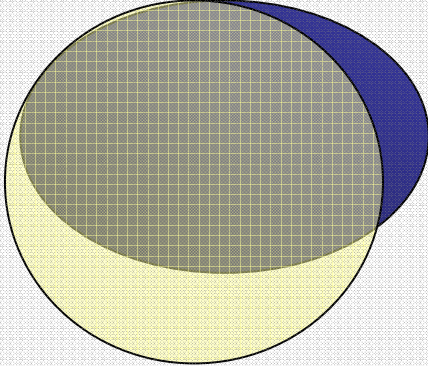
Equilibrium variational principles are intimately related to the structure of equilibrium fluctuations

- Is there a **nonequilibrium** analogy of thermodynamical **fluctuation theory**?

From principles to fluctuation laws

Equilibrium fluctuations

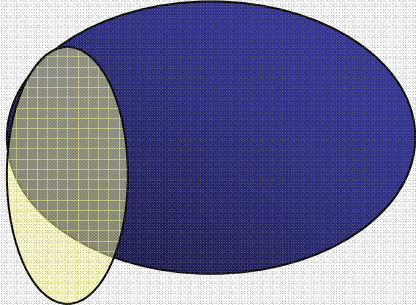
$H(x) = Ne$



$M(x) = Nm_{eq}(e)$
Typical value

The diagram shows a sphere with a grid pattern on its upper half and a yellow lower half. A blue shaded area is on the right side of the sphere.

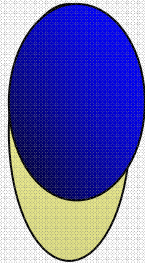
$H(x) = Ne$



$P(M(x) = Nm) = e^{N[s(e,m) - s_{eq}(e)]}$
Probability of fluctuation

The diagram shows a sphere tilted to the right, with a grid pattern on its upper-left part and a yellow lower-left part. A blue shaded area covers the right and top-right parts of the sphere.

add field

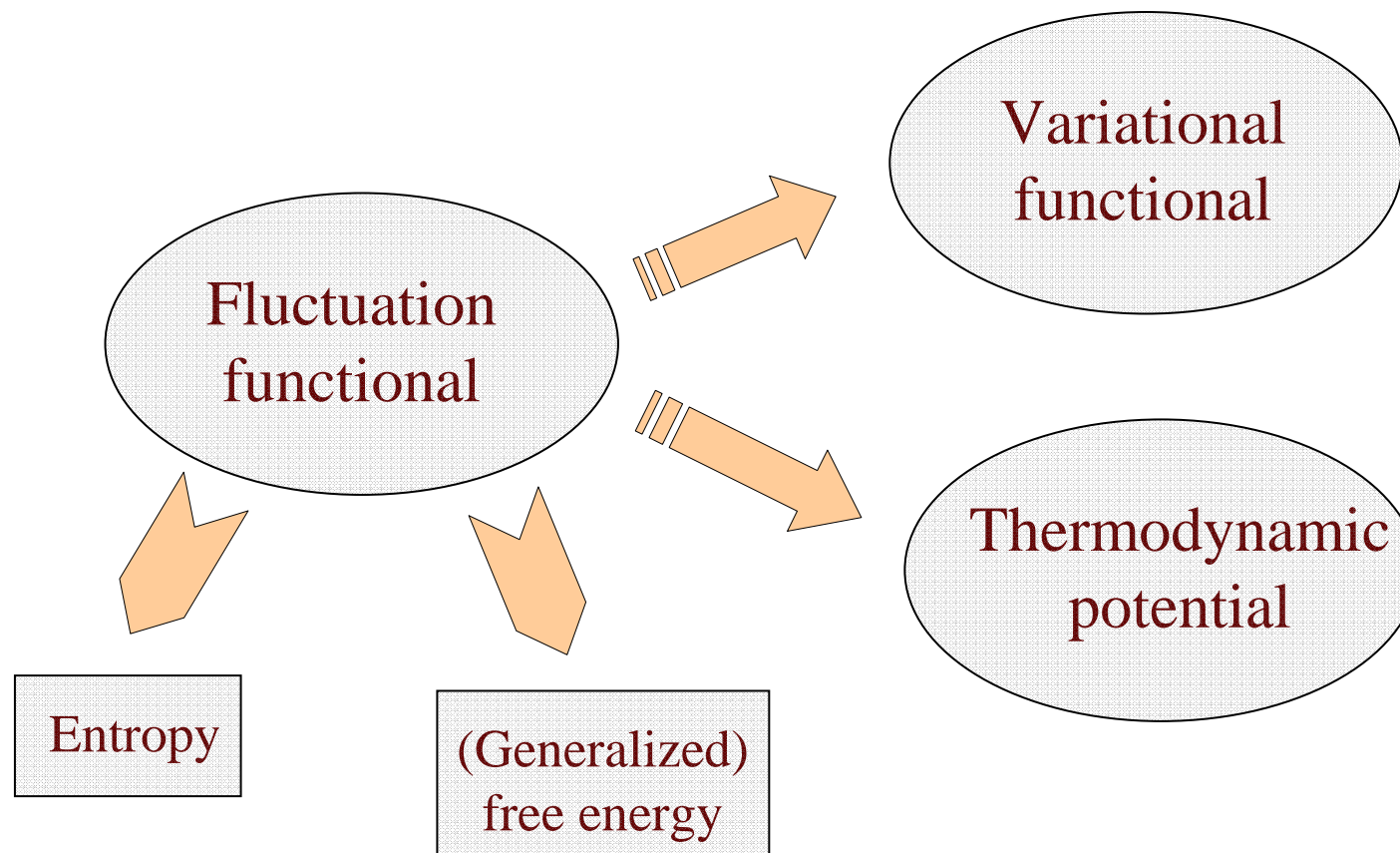


$H^h(x) = H(x) - hM(x) = N[e - hm]$
The fluctuation made typical!
 $s(e, m) = s_{eq}^h(e - hm)$

The diagram shows a sphere tilted to the right, with a blue upper part and a yellow lower part.

From principles to fluctuation laws

Equilibrium fluctuations



From principles to fluctuation laws

Static versus dynamical fluctuations

- Empirical time average:

$$\bar{m}_T = \frac{1}{T} \int_0^T m(x_t) dt$$

- Ergodic property:

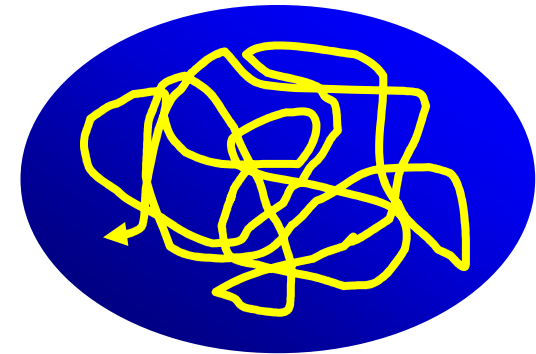
$$\bar{m}_T \rightarrow m_{eq}(e), \quad T \rightarrow \infty$$

- Dynamical fluctuations:

$$P(\bar{m}_T = m) = e^{-T I(m)}$$

- Interpolating between **static** and **dynamical** fluctuations:

$$P\left(\frac{1}{n} \sum_{k=1}^n m(x_{\tau k}) = m\right) = e^{-n I^{(\tau)}(m)}$$



$$H(x) = Ne$$

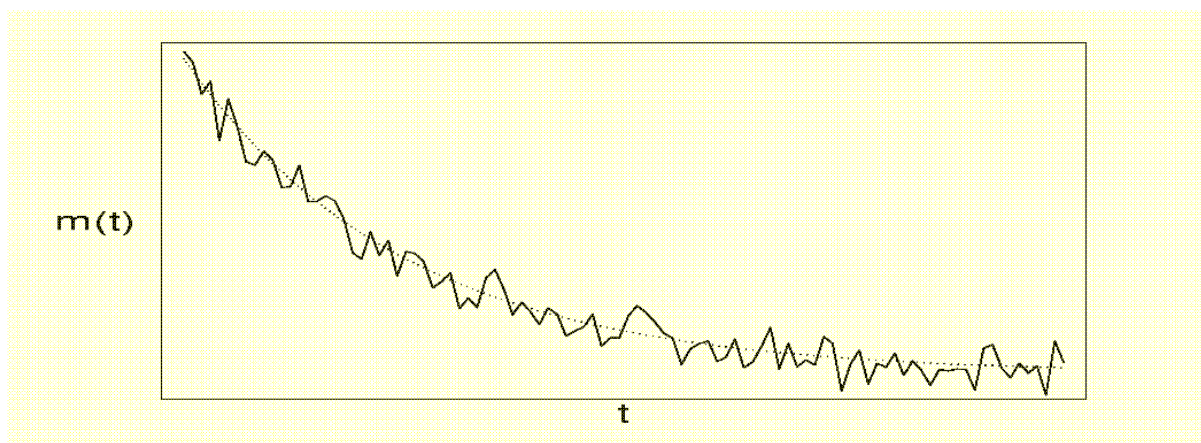
Static: $\tau \rightarrow \infty$

$$I^{(\infty)}(m) = s(e) - s(e, m)$$

Dynamic: $\tau \rightarrow 0$

Effective model of macrofluctuations

Onsager-Machlup theory



- Dynamics: $R dm_t = -s m_t dt + \sqrt{\frac{2R}{N}} dB_t$
- Equilibrium: $P(m_\infty = m) \propto e^{-\frac{1}{2} N s m^2}$
- Path distribution: $P(\omega) = \exp\left[-\frac{N}{4} \int_0^T \frac{R}{2} \left(\frac{dm_t}{dt} + \frac{s}{R} m_t\right)^2\right]$

Effective model of macrofluctuations

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$$P(\omega) = \exp \left[-\frac{N}{4} \int_0^T \frac{R}{2} \left(\frac{dm_t}{dt} + \frac{s}{R} m_t \right)^2 \right]$$

- Dynamical fluctuations:

$$P(\bar{m}_T = m) = P(m_t = m; 0 \leq t \leq T) = \exp \left[-T \frac{Ns^2}{8R} m^2 \right]$$

- (Typical immediate) entropy production rate:

$$\sigma(m) = \frac{dS(m_t)}{dt} = \frac{Ns^2}{2R} m^2$$

Effective model of macrofluctuations

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- Dynamics: $R dm_t = -sm_t dt + \sqrt{\frac{2R}{N}} dB_t$

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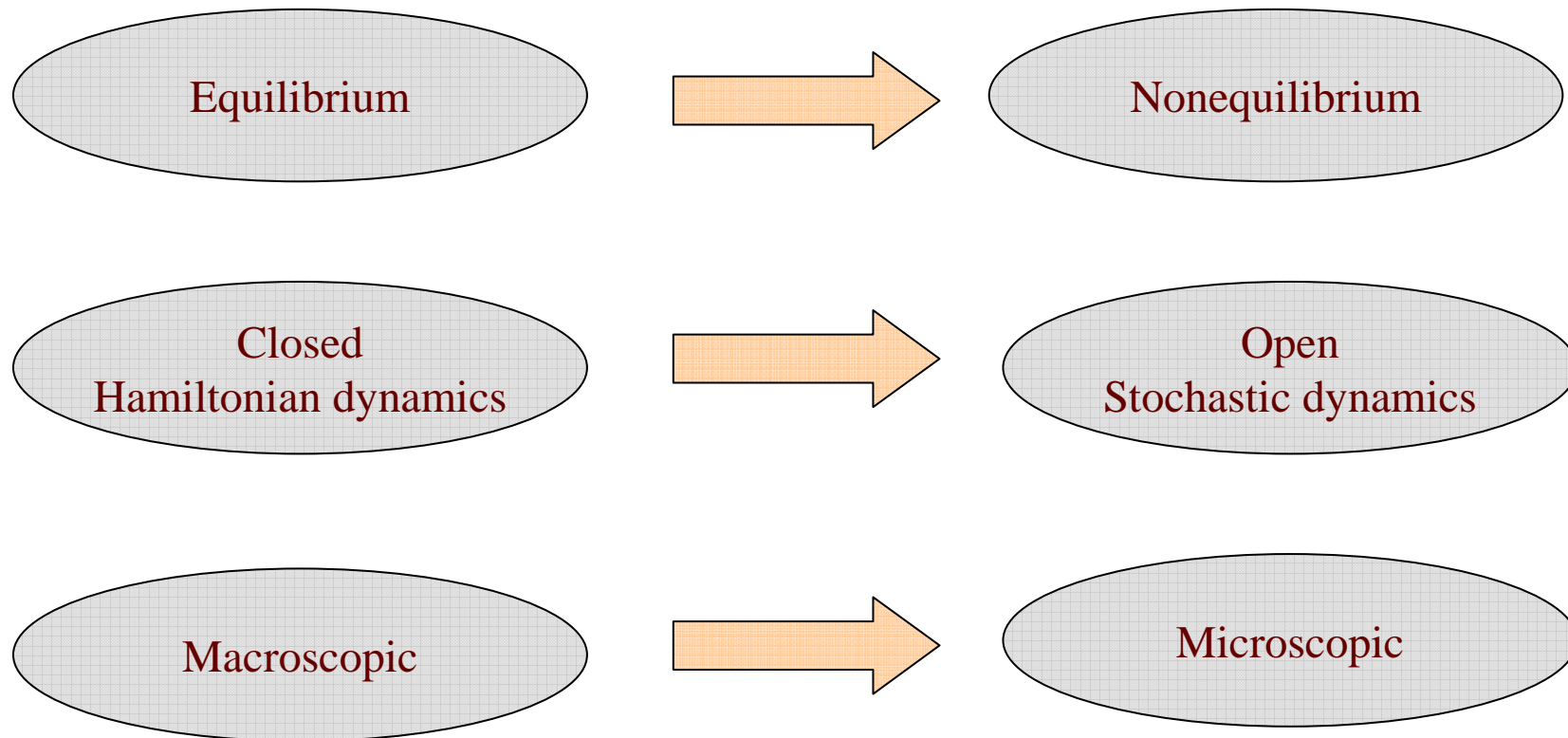
$$P(\bar{m}_T = m) = P(m_t = m; 0 \leq t \leq T) = \exp \left[-T \frac{N s^2}{8R} m^2 \right]$$

$$I(m) = \frac{1}{4} \sigma(m)$$

- (Typical immediate) entropy production rate:

$$\sigma(m) = \frac{dS(m_t)}{dt} = \frac{N s^2}{2R} m^2$$


Towards general theory



Linear electrical networks revisited

Dynamical fluctuations

- Fluctuating dynamics:

$$E = U + R_2 J + E_2^f$$

$$J = C\dot{U} + \frac{U - E_1^f}{R_1}$$

- Johnson-Nyquist noise:

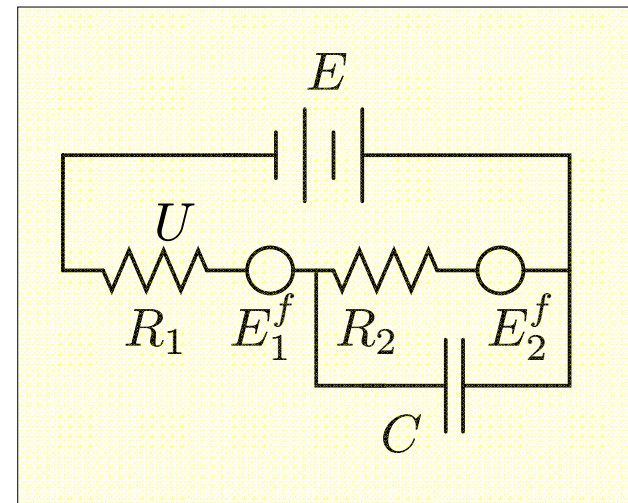
$$E_t^f = \sqrt{\frac{2R}{\beta}} \xi_t \quad \leftarrow \text{white noise}$$

- Empirical time average:

$$\bar{U}_T = \frac{1}{T} \int_0^T U_t dt$$

- Dynamical fluctuation law:

$$-\frac{1}{T} \log P(\bar{U}_T = U) = \frac{1}{4} \frac{\beta_1 \beta_2 (R_1 + R_2)}{\beta_1 R_1 + \beta_2 R_2} \left[\frac{U^2}{R_1} + \frac{(E - U)^2}{R_2} - \frac{E^2}{R_1 + R_2} \right]$$



Linear electrical networks revisited

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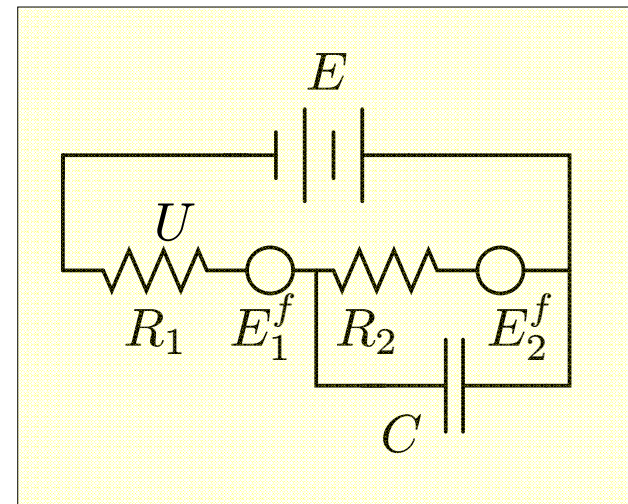
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total dissipated heat

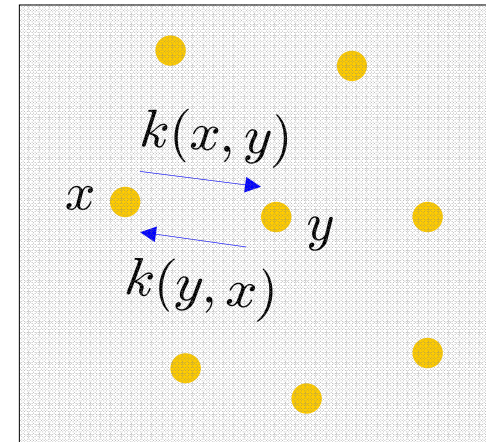


Stochastic models of nonequilibrium

breaking detailed balance

- **Local detailed balance:**

$$\log \frac{k(x,y)}{k(y,x)} = \Delta s(x,y) = -\Delta s(y,x)$$



- **Global detailed balance generally broken:**

$$\Delta s(x,y) = s(y) - s(x) + \epsilon F(x,y)$$

- **Markov dynamics:**


$$\frac{d\rho_t(x)}{dt} = \sum_y [\rho_t(y)k(y,x) - \rho_t(x)k(x,y)]$$

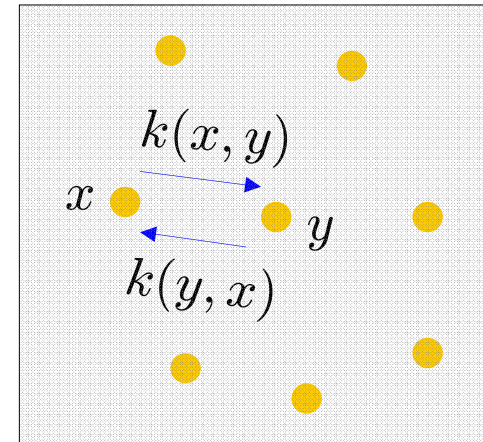
Stochastic models of nonequilibrium

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- Local detailed balance:

$$\log \frac{k(x,y)}{k(y,x)} = \Delta s(x,y) = -\Delta s(y,x)$$

 entropy change
in the environment



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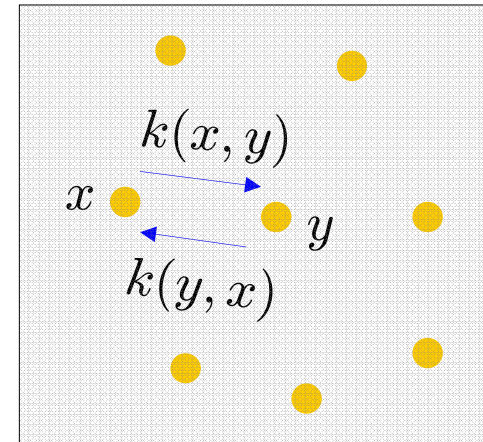
Stochastic models of nonequilibrium

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breaking term

- Markov dynamics:

$$\frac{d\rho_t(x)}{dt} = \sum_y [\rho_t(y)k(y,x) - \rho_t(x)k(x,y)]$$

Stochastic models of nonequilibrium

entropy production

- **Entropy** of the system:

$$S(\rho) = - \sum_x \rho(x) \log \rho(x)$$

- **Mean currents**:

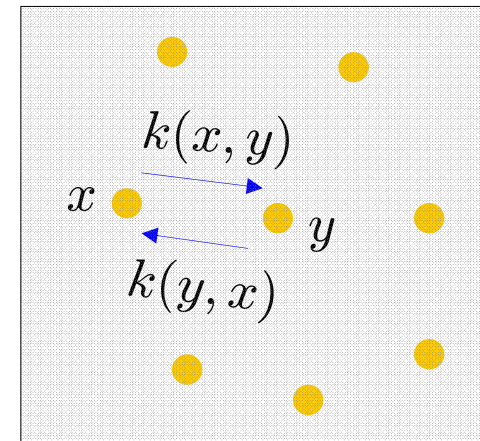
$$J_\rho(x, y) = \rho(x)k(x, y) - \rho(y)k(y, x)$$

zero at detailed balance

- **Mean entropy production rate**:

$$\sigma(\rho) = \frac{dS(\rho_t)}{dt} + \frac{1}{2} \sum_{(x,y)} J_\rho(x, y) \Delta s(x, y)$$

$$= \sum_{x,y} \rho(x)k(x, y) \log \frac{\rho(x)k(x, y)}{\rho(y)k(y, x)}$$



Stochastic models of nonequilibrium

entropy production

- Entropy of the system:

$$S(\rho) = - \sum_x \rho(x) \log \rho(x)$$

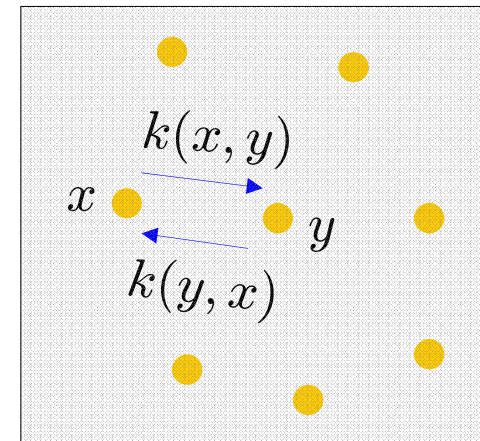
- Entropy fluxes:

$$J_\rho(x, y) = \underbrace{\rho(x)k(x, y) - \rho(y)k(y, x)}$$

zero at detailed balance

- Mean entropy production rate:

$$\begin{aligned} \sigma(\rho) &= \frac{dS(\rho_t)}{dt} + \frac{1}{2} \sum_{(x,y)} J_\rho(x, y) \Delta s(x, y) \\ &= \sum_{x,y} \rho(x)k(x, y) \log \frac{\rho(x)k(x, y)}{\rho(y)k(y, x)} \geq 0 \end{aligned}$$



Warning:
Only for time-reversal
symmetric observables!

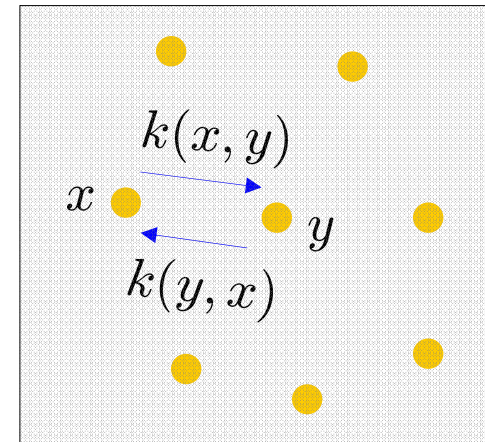
Stochastic models of nonequilibrium

MinEP principle

- (“Microscopic”) **MinEP** principle:

In the first order approximation around detailed balance

$$\sigma(\rho) = \min \Rightarrow \rho = \rho_s + O(\epsilon^2)$$



- Can we again recognize entropy production as a **fluctuation functional**?

Stochastic models of nonequilibrium

dynamical fluctuations

- Empirical occupation times:

$$\bar{p}_T(x) = \frac{1}{T} \int_0^T \chi(\omega_t = x) dt$$

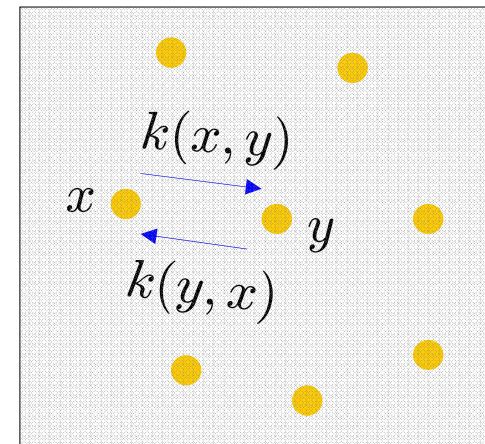
- Ergodic theorem:

$$\bar{p}_T(x) \rightarrow \rho_s(x), \quad T \rightarrow \infty$$

- Fluctuation law for occupation times?

$$P(\bar{p}_T = p) = e^{-T I(p)}$$

- Note: $I(\rho_s) = 0$



Stochastic models of nonequilibrium

dynamical fluctuations

- *Idea*: Make the empirical distribution **typical** by **modifying dynamics**:

$$k(x, y) \longrightarrow k_v(x, y) = k(x, y) e^{[v(y) - v(x)]/2}$$

- The “field” v is such that distribution p is **stationary** distribution for the modified dynamics:

$$\sum_y [p(y)k_v(y, x) - p(x)k_v(x, y)] = 0$$

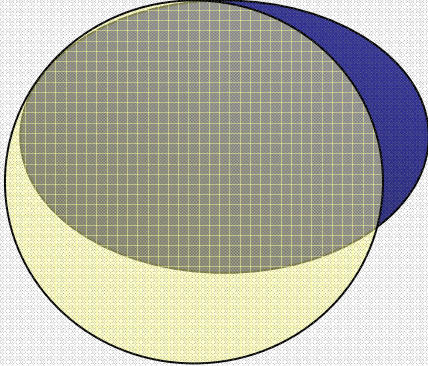
- Comparing both processes yields the **fluctuation law**:

$$I(p) = \sum_{x,y} p(x) [k(x, y) - k_v(x, y)]$$

Recall

Equilibrium fluctuations

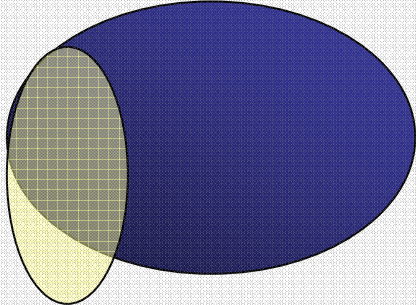
$H(x) = Ne$



$M(x) = Nm_{eq}(e)$
Typical value

The diagram shows a sphere with a grid pattern on its top half and a yellow bottom half. A blue shaded area is on the right side of the sphere.

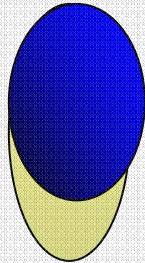
$H(x) = Ne$



$P(M(x) = Nm) = e^{N[s(e,m) - s_{eq}(e)]}$
Probability of fluctuation

The diagram shows a sphere tilted to the right, with a grid pattern on its top-left side and a yellow bottom-left side. The rest of the sphere is blue.

add field



$H^h(x) = H(x) - hM(x) = N[e - hm]$
The fluctuation made typical!
 $s(e, m) = s_{eq}^h(e - hm)$

The diagram shows a sphere tilted to the right, with a blue top half and a yellow bottom half.

Stochastic models of nonequilibrium

dynamical fluctuations

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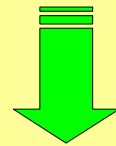
$$I(p) = \sum_{x, y} p(x) [k(x, y) - k_v(x, y)]$$

Stochastic models of nonequilibrium

dynamical

Traffic = mean dynamical activity:

$$\mathcal{T} = \frac{1}{2} \sum_{x,y} p(x)k(x,y) + p(y)k(y,x)$$



$I(p)$ = excess in traffic

- Ideal
model

$k(x,y)$

- The
distribution

\sum_y

- Comparing both processes yields the **fluctuation law:**

$$I(p) = \sum_{x,y} p(x) [k(x,y) - k_v(x,y)]$$

Stochastic models of nonequilibrium

Recall: entropy production functional

- **Entropy** of the system:

$$S(\rho) = - \sum_x \rho(x) \log \rho(x)$$

- **Mean currents**:

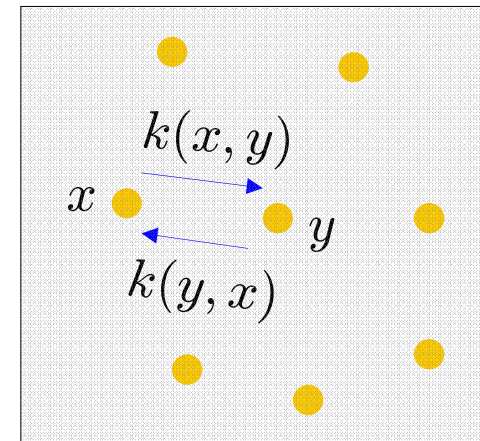
$$J_\rho(x, y) = \rho(x)k(x, y) - \rho(y)k(y, x)$$

zero at detailed balance

- **Mean entropy production rate**:

$$\sigma(\rho) = \frac{dS(\rho_t)}{dt} + \frac{1}{2} \sum_{(x,y)} J_\rho(x, y) \Delta s(x, y)$$

$$= \sum_{x,y} \rho(x)k(x, y) \log \frac{\rho(x)k(x, y)}{\rho(y)k(y, x)}$$



Stochastic models of nonequilibrium

dynamical fluctuations close to equilibrium

- General observation:

 *In the first order approximation around detailed balance*

$$I(p) = \frac{1}{4} [\sigma(p) - \sigma(\rho_s)] + o(\epsilon^2)$$

- The variational functional is recognized as an **approximate** fluctuation functional
- *A consequence:* A natural way how to go **beyond MinEP principle** is to systematically analyze appropriate fluctuation laws

Stochastic models of nonequilibrium

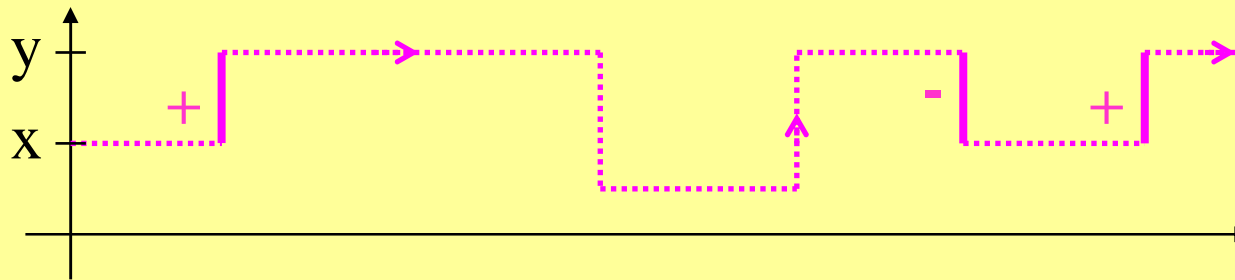
dynamical fluctuations close to equilibrium

- General observation:

In the first order approximation around detailed balance

Empirical currents:

$$\bar{J}_T(x, y) = \frac{1}{T} [\#\{\text{jumps } x \rightarrow y \text{ in } [0, T]\} - \#\{\text{jumps } y \rightarrow x \text{ in } [0, T]\}]$$



an

beyond
tion laws

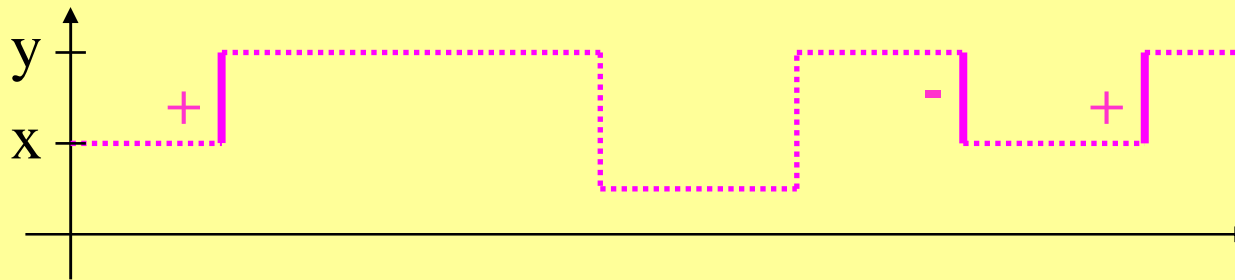
Stochastic model of dynamical fluctuations

- General observations

In the first part, we will be more detailed

Empirical currents:

$$\bar{J}_T(x, y) = \frac{1}{T} [\#\{\text{jumps } x \rightarrow y\} - \#\{\text{jumps } y \rightarrow x\}]$$



Typically,

$$\bar{J}_T(x, y) \rightarrow \rho_s(x)k(x, y) - \rho_s(y)k(y, x)$$

Fluctuation law: $J_s(x, y)$

$$P(\bar{J}_T = J) = e^{-T G(J)}$$

with the fluctuation functional

$$G(J) = \frac{1}{4} [\dot{S}(J_s) - \dot{S}(J)] + o(\epsilon^2)$$

on stationary currents satisfying

$$\dot{S}(J) = D(J)$$

beyond fluctuation laws

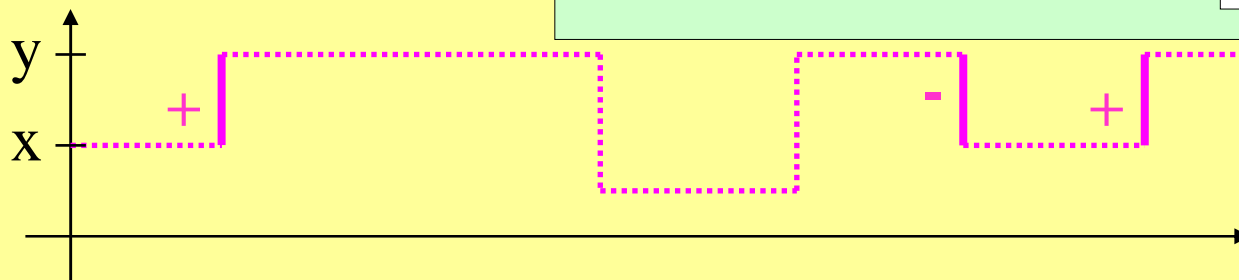
Stochastic model dynamical fluctuations

- General observations

*In the first part
detailed*

Empirical currents:

$$\bar{J}_T(x, y) = \frac{1}{T} [\#\{\text{jumps } x \rightarrow y\} - \#\{\text{jumps } y \rightarrow x\}]$$



Typically,

$$\bar{J}_T(x, y) \rightarrow \rho_s(x)k(x, y) - \rho_s(y)k(y, x)$$

Fluctuation law:

$$P(\bar{J}_T = J) = e^{-T G(J)}$$

Entropy flux
 $\frac{1}{2} \sum_{x,y} J(x, y) \Delta s(x, y)$

with the fluctuation functional

$$G(J) = \frac{1}{4} [\dot{S}(J_s) - \dot{S}(J)] + o(\epsilon^2)$$

on stationary currents satisfying

$$\dot{S}(J) = D(J)$$

Onsager dissipation
function

beyond
fluctuation laws

Stochastic models of nonequilibrium

towards general fluctuation theory

- It is useful to study the **occupation time** statistics and **current** statistics **jointly**
- Joint occupation-current statistics has a **canonical structure**

Driving-parameterized dynamics

$$k_F(x, y) = k_0(x, y) e^{F(x, y)/2}$$

Reference equilibrium

anti-symmetric

Current potential function

$$H(p, F) = 2[\mathcal{T}_F(p) - \mathcal{T}_0(p)]$$

Traffic

Canonical equations

$$\left. \frac{\delta H}{\delta F(x,y)} \right|_{p,F} = J_F(x,y) \quad \overset{\text{Legendre}}{\longleftrightarrow} \quad \left. \frac{\delta G}{\delta J(x,y)} \right|_{p,J_F} = F(x,y)$$

Joint occupation-current fluctuation functional

$$\mathcal{I}_F(p, J) = \frac{1}{2} [G(p, J) + H(p, F) - \dot{S}(F, J)]$$

Driving-parameterized dynamics

$$k_F(x, y) = k_0(x, y) e^{F(x,y)/2}$$

Reference equilibrium

anti-symmetric

Current potential function

$$H(p, F) = 2[\mathcal{T}_F(p) - \mathcal{T}_0(p)]$$

Traffic

Stochastic models of nonequilibrium

consequences of canonical formalism

- Functional G describes (reference) **equilibrium** dynamical fluctuations
- **Fluctuation symmetry** immediately follows:
$$\mathcal{I}_F(p, -J) - \mathcal{I}_F(p, J) = \dot{S}(F, J)$$
- **Symmetric** (p) and **antisymmetric** (J) fluctuations are **coupled** away from equilibrium, but:

Decoupling between p and J

- for small fluctuations
- close to equilibrium

General conclusions

what we know

- Both **MinEP** and **MaxEP** principles naturally follow from the fluctuation laws for **empirical occupation times** and **empirical currents**, respectively
- The validity of both principles is restricted to the **close-to-equilibrium** regime and it is essentially a consequence of
 - **decoupling** between time-symmetric and time-antisymmetric fluctuations
 - intimate relation between **traffic** and **entropy production** for Markovian dynamics close to detailed balance
- **Time-symmetric** fluctuations are in general governed by the **traffic** functional (nonperturbative result!)
- **Joint** occupation-current fluctuations have a general **canonical structure**, generalizing the original Onsager-Machlup theory
- Our approach can be extended to **semi-Markov** systems with some similar conclusions, cf. [6]



General conclusions

what we would like to know

- What is the **operational meaning** of new quantities (traffic,...) emerging in the dynamical fluctuation theory?
- Are there useful **computational schemes** for the fluctuation functionals and can one **systematically improve** on the EP principles beyond equilibrium?
- What is the relation between **static** and **dynamical** fluctuations?
- Could the dynamical fluctuation theory be a useful approach towards building **nonequilibrium thermodynamics beyond close-to-equilibrium**?

...and still many other things would be nice to know...



References

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- 6) C. Maes, K.N., and B. Wynants, *J. Phys. A: Math. Theor.* **42**, 365002 (2009)

A high-altitude mountain landscape featuring snow-covered peaks and a glacier. The sky is a clear, deep blue. The foreground shows a rocky, snow-dusted slope leading down to a large, textured glacier. The background consists of several jagged, snow-capped mountain peaks under a clear blue sky.

*Thank You
for Your Attention!*