

Cluster expansion using Penrose tree-graph identity

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Joint work with Aldo Procacci

Warwick Statistical Mechanics Seminar

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- Cluster expansion method in that case (dilute gas).
- Generalization to polymers.

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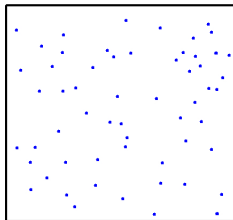
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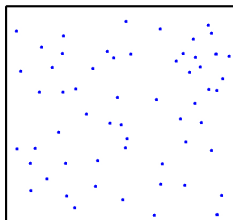


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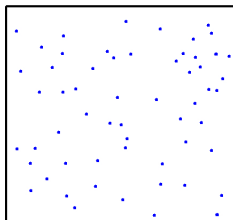
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divided by
$$Z(\Lambda) = \sum_{m \geq 0} \int_{\Lambda^m} d\vec{y} \frac{z^m}{m!} e^{-\beta \sum_{i,j} V(y_i - y_j)}$$

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$\beta > 0$, inverse temperature.

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If $V \equiv 0$, it is Poisson point process, a.k.a. *ideal classical gas*.

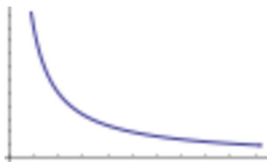
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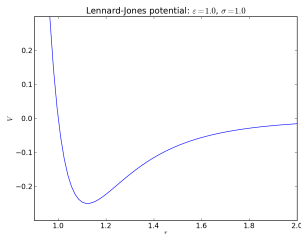
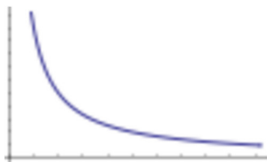
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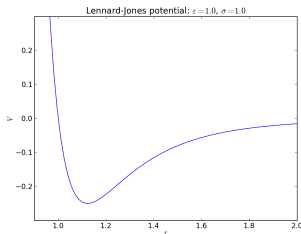
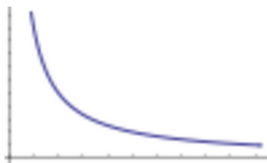
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Low interactions: β or z small

Partition function:

$$\Lambda \subset \mathbb{R}^d, \quad Z(\Lambda) = \sum_{m \geq 0} \frac{z^m}{m!} \int_{\Lambda^m} d\vec{y} \, e^{-\beta \sum_{i,j} V(y_i - y_j)}$$

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Uniform convergence radius for $\Lambda \rightarrow \mathbb{R}^d$?

1963: Penrose, Ruelle First positive answer.

Later: Better bounds but more hypothesis, using trees & graphs.

Here: Better bounds, no extra hypothesis.

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Because $\alpha(g_1)\alpha(g_2) = \alpha(g_1 \cup g_2)$ and

the series of Exp combines connected components in every possible way.

We need to bound

$$\left| \sum_{g \in \mathcal{C}_X} \prod_{s \in E_g} (e^{-\beta V(s)} - 1) \right|$$

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proof:

$$\prod_{s \in M(\tau) \setminus \tau} ((e^{-V_s} - 1) + 1) = \sum_{g' \subset M(\tau) \setminus \tau} \left(\prod_{s \in g'} (e^{-V_s} - 1) \right)$$

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Proposition: $T^{-1}(\tau) = [\tau, M(\tau)]$

where M is easy to describe.

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Kruskal algorithm: greedy construction of minimum spanning tree.

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If $\tau \subset g \subset M(\tau)$ then $T(g) = \tau$.

Penrose tree-graph identity for minimum spanning trees wrt V .

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Further trick

$$= \sum_{\tau \in \mathcal{T}_X} \prod_{s \in \tau} (1 - e^{-|V_s|}) \prod_{s \in M(\tau) \setminus \tau^+} e^{-V_s}$$

Stability hypothesis

the potential V has to satisfy

$$\sum_{\{x,y\} \in X^{(2)}} V(x-y) \geq -B|X|$$

for any finite configuration $X \subset \mathbb{R}^d$, for some constant $B \geq 0$.

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If $V \geq 0$ ✓

If not, negative part has to be controlled.

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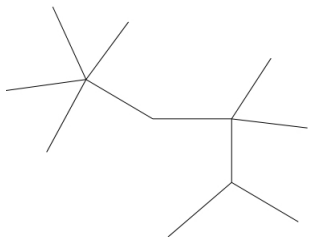
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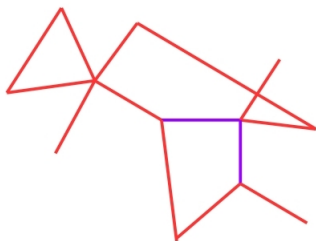
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(as if the graph $M(\tau) \setminus \tau^+$ was the complete graph)

Tree: τ



$$\sum_{s \in M(\tau) \setminus \tau^+} V_s \geq$$



$$\geq \sum_{cc} \sum_s V_s \geq -B|X|$$



$$\leq \sum_{\tau \in \mathcal{T}_X} \prod_{s \in \tau} (1 - e^{-|V_s|}) e^{B|X|}$$

integrate over configurations of points, X

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For $|X| = n$,

$$\int (...) \leq z^n \frac{n^{n-2}}{n!} \left(\int_{\mathbb{R}^d} (1 - e^{-\beta|V(x)|}) dx \right)^{n-1} e^{\beta B n |D|}$$

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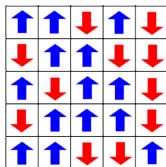
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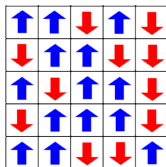
$$R \geq \left(e^{\beta B + 1} \int_{\mathbb{R}^d} (1 - e^{-\beta|V|}) \right)^{-1}$$

Ising Model



Energy of a configuration:
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$$p_{\Lambda}(\sigma) = \frac{e^{-\beta H(\sigma)}}{\sum_{\sigma'} e^{-\beta H(\sigma')}} \quad (\text{Boltzmann-Gibbs})$$

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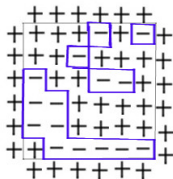
For β small, weak interactions, there is only one Gibbs measure.

$d \geq 2$, for β large, strong interactions, at least 2 Gibbs measures.

(so if $d \geq 2$, we have a *phase transition*)

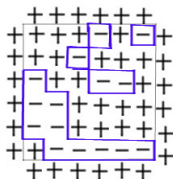
With cluster expansion we can understand both statements, and obtain explicit formulas.

For β large:



Ising configuration = polymer configuration.

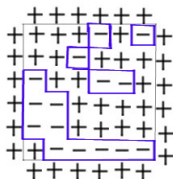
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$$p(\sigma) = \frac{\prod_{\gamma \in \sigma} e^{-\beta|\gamma|}}{\sum_{\substack{r \in \mathbb{N}_0 \\ (\gamma_1, \dots, \gamma_r)}} \frac{1}{r!} \prod_{i=1}^r e^{-\beta|\gamma_i|}}$$

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Polymer activity: $e^{-\beta \text{ length}}$

Interaction: contact forbidden.

β small: polymers are defined in a different way.

$$Z(\Lambda) = \int_{\Gamma_\Lambda} dX \prod_{\{\gamma_1, \gamma_2\}} e^{-V(\gamma_1, \gamma_2)}$$

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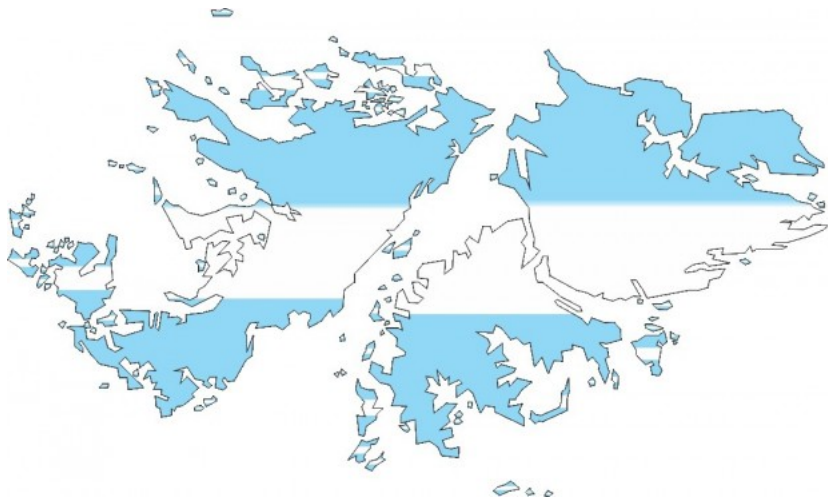
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The same method works



Thank you!